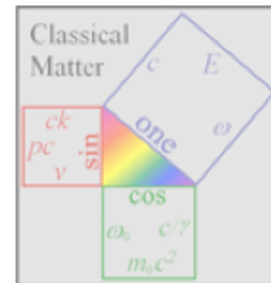


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Dirac Equation for Spin Density in an Ideal Elastic Solid

Robert Close, PhD.

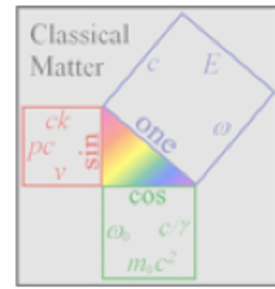
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Spin Density



Belinfante-Rosenfeld symmetric stress-energy tensor implies dynamic (wave) & conjugate momenta:

$$\mathbf{P}_{Total} = -\text{Re}[\psi^\dagger i\nabla\psi] + \frac{1}{2}\nabla \times \left[\psi^\dagger \frac{\boldsymbol{\sigma}}{2} \psi \right]$$

$$= \mathbf{P}_{wave} + \frac{1}{2}\nabla \times \mathbf{s}$$

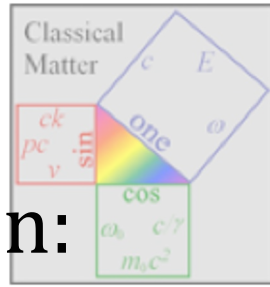
$$\mathbf{J} = -\text{Re}(\mathbf{r} \times [\psi^\dagger i\nabla\psi]) + \left[\psi^\dagger \frac{\boldsymbol{\sigma}}{2} \psi \right]$$

$$= \mathbf{L}_{wave} + \mathbf{s}$$

Elastic waves in a solid also have dynamic & conjugate (or “intrinsic”) momenta.

Is Standard Model a decomposition of elastic waves into “particles”?

Classical Spin Density



Incompressible Helmholtz decomposition:

$$\mathbf{p} \equiv \rho \mathbf{u} = \mathbf{p}_0 + \nabla \Phi + \frac{1}{2} \nabla \times \mathbf{s}$$

Angular momentum:

$$\mathbf{S} = \int \mathbf{r} \times \mathbf{p} d^3 r = \int \mathbf{s} d^3 r + \text{b.t.}$$

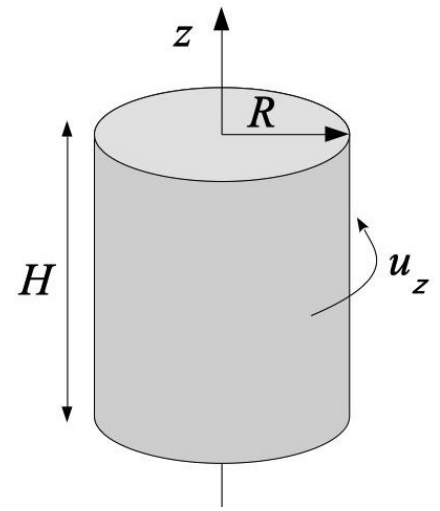
Kinetic energy: $(\mathbf{w} = \frac{1}{2} \nabla \times \mathbf{u})$

$$K = \int \frac{1}{2} \rho u^2 d^3 r = \int \frac{1}{2} \mathbf{w} \cdot \mathbf{s} d^3 r + \text{b.t.}$$

Spin density is the momentum conjugate to angular velocity ($\mathcal{L} \sim K; \frac{\partial \mathcal{L}}{\partial \mathbf{w}} = \frac{\mathbf{s}}{2} + \frac{\mathbf{s}}{2} = \mathbf{s}$).

Rotating Cylinder

Parabolic radial distribution of spin density.



$$s_z = \rho w_0 [R^2 - r^2]$$

$$r \leq R$$

$$u_\phi = \frac{1}{2\rho} \frac{\partial}{\partial r} s_z = r w_0$$

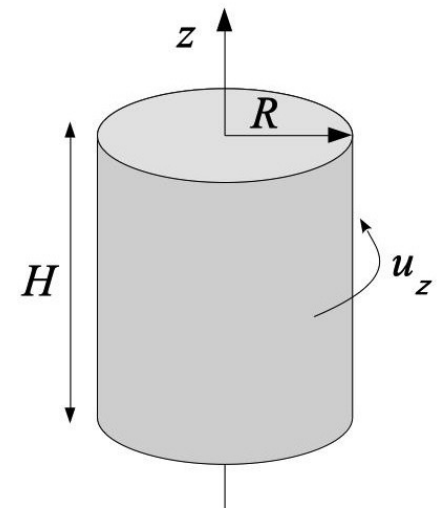
$$r \leq R$$

$$w_z = \frac{1}{2r} \frac{\partial}{\partial r} (r u_\phi)$$

$$= w_0 [1 - R\delta(R - r)/2] \quad r \leq R$$

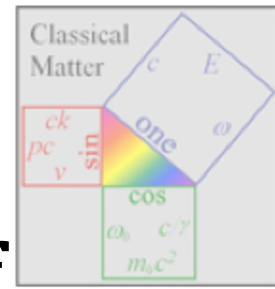
Rotating Cylinder

Integration yields the usual total angular momentum and kinetic energy.



$$\begin{aligned} S_z &= \int \rho w_0 [R^2 - r^2] d^3 r \\ &= 2\pi\rho H \left(\frac{R^4}{4} \right) w_0 = \frac{MR^2}{2} w_0 = I w_0 \\ K &= \int \frac{1}{2} \mathbf{w} \cdot \mathbf{s} d^3 r = \frac{1}{2} I w_0^2 \end{aligned}$$

Equation of Evolution



Momentum density: $\partial_t \mathbf{p} + \mathbf{u} \cdot \nabla \mathbf{p} = \mathbf{f}$

Changes due to convection & force.

Spin density: $\partial_t \mathbf{s} + \mathbf{u} \cdot \nabla \mathbf{s} - \mathbf{w} \times \mathbf{s} = \boldsymbol{\tau}$

Changes \sim convection, rotation, & torque.

Let $\mathbf{s} \equiv \partial_t \mathbf{Q}$, Displacement: $\boldsymbol{\xi} = \frac{1}{2\rho} \nabla \times \mathbf{Q}$

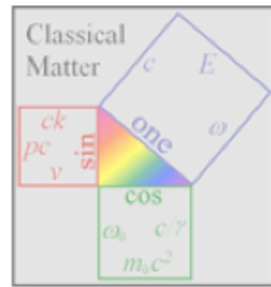
Elastic solid: Torque = $\boldsymbol{\tau} = -\frac{\mu}{\rho} \nabla \times \nabla \times \mathbf{Q}$

$$\partial_t^2 \mathbf{Q} + \mathbf{u} \cdot \nabla \partial_t \mathbf{Q} - \mathbf{w} \times \partial_t \mathbf{Q} = c^2 \nabla^2 \mathbf{Q}$$

Nonlinearity \Rightarrow soliton solutions

Factor the Wave Equation

$$Q(z, t) = Q_F(ct - z) + Q_B(ct + z)$$



- Independent states 180° apart \Rightarrow spin $1/2$
- The second-order wave equations are:

$$\partial_t^2 Q_B - c^2 \partial_z^2 Q_B = 0$$

$$\partial_t^2 Q_F - c^2 \partial_z^2 Q_F = 0$$

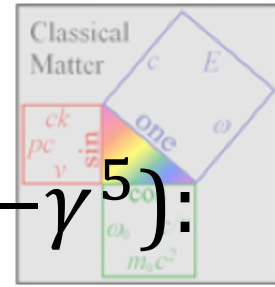
- Equivalent to first-order matrix equation:

$$\partial_t \begin{bmatrix} \dot{Q}_B \\ \dot{Q}_F \end{bmatrix} - c \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \partial_z \begin{bmatrix} \dot{Q}_B \\ \dot{Q}_F \end{bmatrix} = 0$$

- The Dirac equation of quantum mechanics simply extends this matrix equation to vectors in 3D.

Factor the Wave Equation

First-order 1-D wave eqn. (chiral $\beta^3 = -\gamma^5$):



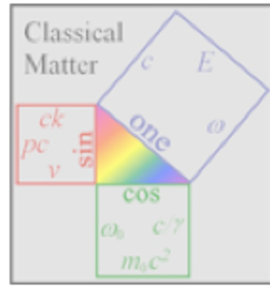
$$\partial_t(\psi^T \sigma_z \psi) - c \partial_z(\psi^T \beta^3 \psi) = 0$$

$$\partial_t Q = \begin{bmatrix} (\dot{Q}_{B+})^{1/2} \\ (\dot{Q}_{F-})^{1/2} \\ (\dot{Q}_{F+})^{1/2} \\ (\dot{Q}_{B-})^{1/2} \end{bmatrix}^T \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{bmatrix} (\dot{Q}_{B+})^{1/2} \\ (\dot{Q}_{F-})^{1/2} \\ (\dot{Q}_{F+})^{1/2} \\ (\dot{Q}_{B-})^{1/2} \end{bmatrix} = \psi^T \sigma_z \psi$$

$$c \partial_z Q = \begin{bmatrix} (\dot{Q}_{B+})^{1/2} \\ (\dot{Q}_{F-})^{1/2} \\ (\dot{Q}_{F+})^{1/2} \\ (\dot{Q}_{B-})^{1/2} \end{bmatrix}^T \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{bmatrix} (\dot{Q}_{B+})^{1/2} \\ (\dot{Q}_{F-})^{1/2} \\ (\dot{Q}_{F+})^{1/2} \\ (\dot{Q}_{B-})^{1/2} \end{bmatrix} = \psi^T \beta^3 \psi$$

Factor the Wave Equation

Matrices for 3-D polarization & wave velocity:



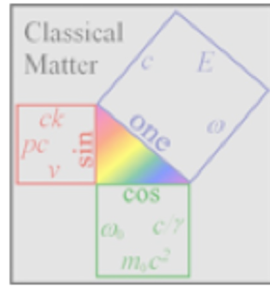
$$\sigma_x = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

$$\sigma_y = \begin{pmatrix} 0 & -i & 0 & 0 \\ i & 0 & 0 & 0 \\ 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \end{pmatrix}$$

$$\sigma_z = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

Factor the Wave Equation

Matrices for 3-D spatial derivatives:



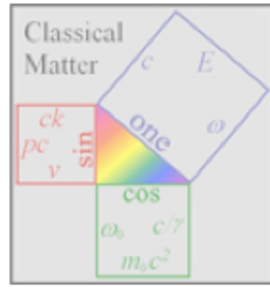
$$\beta^1 = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix} = \gamma^0 \perp \mathbf{s} \rightarrow \text{Curl}$$

$$\beta^2 = \begin{pmatrix} 0 & 0 & -i & 0 \\ 0 & 0 & 0 & -i \\ i & 0 & 0 & 0 \\ 0 & i & 0 & 0 \end{pmatrix} = i\gamma^5\gamma^0 \perp \mathbf{s} \rightarrow 0$$

$$\beta^3 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} = -\gamma^5 \parallel \mathbf{s} \rightarrow \text{Div}$$

Wave velocity matrices: $-c\beta^3\sigma_i = c\gamma^5\sigma_i$

Linear Dirac Equation



1-D wave equation:

$$0 = \psi^T \sigma_z \{ \partial_t \psi + c \gamma^5 \sigma_z \partial_z \psi \} + \text{Transp.}$$

$$0 = \partial_t^2 Q_z - c^2 \partial_z^2 Q_z$$

3-D wave equation (replace z with i or j):

$$\psi^\dagger \sigma_i [\partial_t \psi + c \gamma^5 \sigma_j \partial_j \psi] + adj. = 0$$

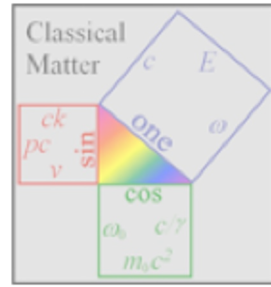
Same as electron! Mass term drops out.

$$\psi^\dagger \sigma_i [\partial_t \psi + c \gamma^5 \sigma_j \partial_j \psi + i M \gamma^0 \psi] + adj. = 0$$

Wave function describes spin density:

$$s_i = (1/2) \psi^\dagger \sigma_i \psi$$

Dirac Equation



Linear vector wave equations:

$$0 = \partial_t [\psi^\dagger \boldsymbol{\sigma} \psi] + c \nabla [\psi^\dagger \gamma^5 \psi] - ic [\nabla \psi^\dagger \times \gamma^5 \boldsymbol{\sigma} \psi + \psi^\dagger \gamma^5 \boldsymbol{\sigma} \times \nabla \psi]$$

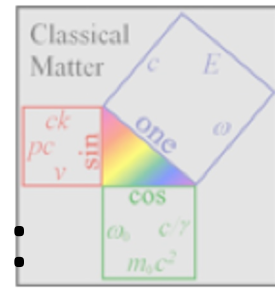
$$0 = \partial_t^2 \mathbf{Q} - c^2 \nabla (\nabla \cdot \mathbf{Q}) + c^2 \nabla \times \nabla \times \mathbf{Q} = \partial_t^2 \mathbf{Q} - c^2 \nabla^2 \mathbf{Q}$$

$$\mathbf{s} = \partial_t \mathbf{Q} = (1/2) \psi^\dagger \boldsymbol{\sigma} \psi$$

$$c \nabla \cdot \mathbf{Q} = -(1/2) \psi^\dagger \gamma^5 \psi$$

$$c^2 \nabla \times \nabla \times \mathbf{Q} = \frac{-ic}{2} [\nabla \psi^\dagger \times \boldsymbol{\sigma} \psi + \psi^\dagger \gamma^5 \boldsymbol{\sigma} \times \nabla \psi]$$

Dirac Equation



Nonlinear Dirac eqn. $\mathbf{s} \equiv \partial_t \mathbf{Q} = \left[\psi^\dagger \frac{\boldsymbol{\sigma}}{2} \psi \right]:$

$\partial_t \psi =$

$$c \beta^3 \boldsymbol{\sigma} \cdot \nabla \psi - \mathbf{u} \cdot \nabla \psi - \frac{i}{2} \tilde{\omega}_0 \gamma^0 \psi - \frac{i}{2} \mathbf{w} \cdot \boldsymbol{\sigma} \psi$$

Propagation, convection, rotation of wave velocity, rotation of the solid medium.

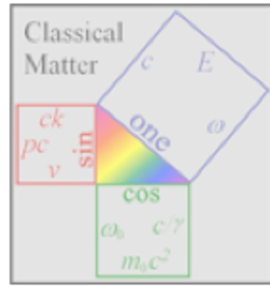
Mass represents rotation of wave velocity.*

Each term has a simple physical interpretation.

For plane waves, velocity rotation cancels rotation of the medium.

*c.f. Hestenes, *Found. Phys.* **20**(10):1213-32,1990

Lagrangian Density



\mathbf{w} & \check{w}_0 terms contribute twice to E-L eqn.

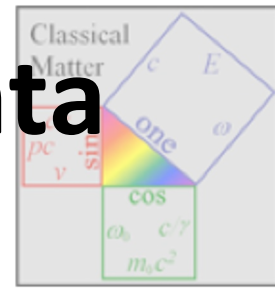
$$\mathcal{L} = \text{Re}\{\psi^\dagger i \partial_t \psi + c\psi^\dagger \gamma^5 \boldsymbol{\sigma} \cdot i\nabla\psi + \mathbf{u} \cdot \psi^\dagger i\nabla\psi\}$$

$$- \frac{1}{4} \check{w}_0 \psi^\dagger \gamma^0 \psi - \frac{1}{4} \mathbf{w} \cdot \psi^\dagger \boldsymbol{\sigma} \psi$$

$$\mathbf{u} = \frac{1}{2\rho} \nabla \times \mathbf{s} = \frac{1}{2\rho} \nabla \times \left(\psi^\dagger \frac{\boldsymbol{\sigma}}{2} \psi \right);$$

$$\mathbf{w} = \frac{1}{2} \nabla \times \mathbf{u} = \frac{1}{4\rho} \nabla \times \nabla \times \left(\psi^\dagger \frac{\boldsymbol{\sigma}}{2} \psi \right)$$

$$\check{w}_0 = \frac{1}{2\rho c} f = \frac{1}{4\rho} \nabla^2 \left(\psi^\dagger \frac{\gamma^0}{2} \psi \right)$$



Dynamic & Conjugate Momenta

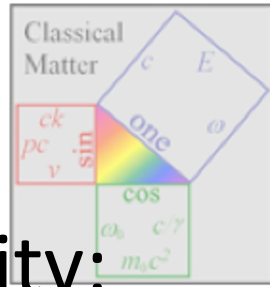
$$\mathbf{P} = -[\psi^\dagger i\nabla\psi] + \frac{1}{2}\nabla \times \left[\psi^\dagger \frac{\boldsymbol{\sigma}}{2} \psi \right]$$

$$= \mathbf{P}_{wave} + \mathbf{p}_{solid}$$

$$\mathbf{J} = -\mathbf{r} \times [\psi^\dagger i\nabla\psi] + \left[\psi^\dagger \frac{\boldsymbol{\sigma}}{2} \psi \right]$$

$$= \mathbf{L}_{wave} + \mathbf{s}_{solid}$$

- Momentum consistent with Belinfante-Rosenfeld energy-momentum tensor.
- Wave (orbital) and spin angular momentum.



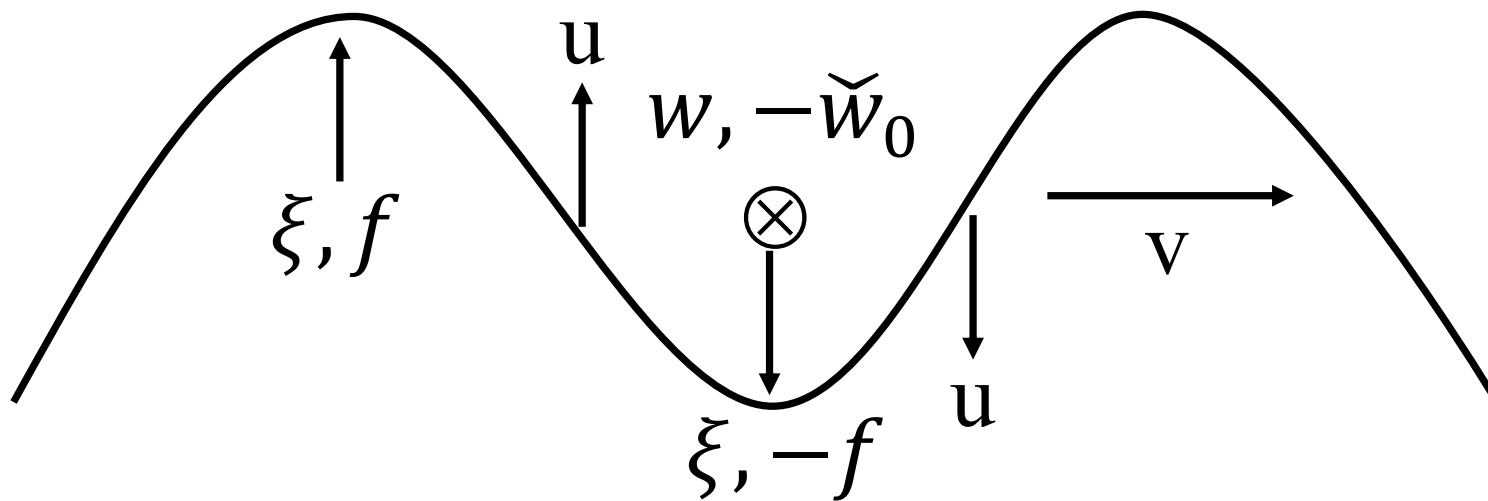
Potential Energy

$\psi^\dagger \gamma^0 \psi \sim$ displacement. $\check{w}_0 \sim$ force density:

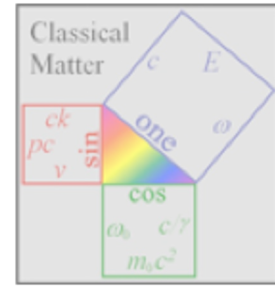
$$\check{w}_0 = \frac{1}{2\rho c} f = \frac{1}{4\rho} \nabla^2 \left(\psi^\dagger \frac{\gamma^0}{2} \psi \right)$$

PE density is: $U = \frac{\check{w}_0}{4} \psi^\dagger \gamma^0 \psi \sim -\frac{1}{2} \mathbf{f} \cdot \boldsymbol{\xi}$

Rotational PE: $U_R = \mathcal{E} - U$



Hamiltonian Density



$$\begin{aligned}
 \mathcal{H} &= \text{Re}(-c\psi^\dagger \gamma^5 \boldsymbol{\sigma} \cdot i\nabla\psi - \mathbf{u} \cdot \psi^\dagger i\nabla\psi) \\
 &\quad + \frac{\check{\omega}_0}{4} \psi^\dagger \gamma^0 \psi + \frac{\mathbf{w}}{4} \cdot \psi^\dagger \boldsymbol{\sigma} \psi \\
 &= \left(\varepsilon - 0 + \frac{1}{2} \mathbf{f} \cdot \boldsymbol{\xi} \right) + \frac{1}{2} \mathbf{w} \cdot \mathbf{s} \\
 &= U_R + K_R
 \end{aligned}$$

Rotational potential energy: $U_R = \varepsilon + \frac{1}{2} \mathbf{f} \cdot \boldsymbol{\xi}$

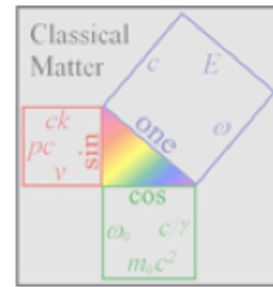
Rotational kinetic energy: $K_R = \frac{1}{2} \mathbf{w} \cdot \mathbf{s} = \varepsilon - \frac{1}{2} \rho u^2$

Equal integrals for vectors & bispinors:

Vectors: Wave velocity \times wave momentum = $2U$

Bispinors: Wave velocity \times wave momentum = ε

Operators as Generators

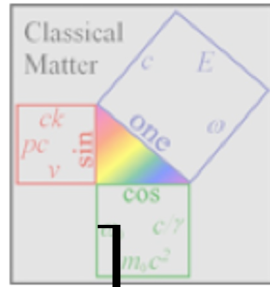


$$\mathbf{P}\psi = -i\nabla\psi + \frac{1}{2}\nabla \times \frac{\boldsymbol{\sigma}}{2}\psi$$

$$\mathbf{J}\psi = -\mathbf{r} \times i\nabla\psi + \frac{\boldsymbol{\sigma}}{2}\psi$$

- \mathbf{J} is generator of rotations (orbital for coordinate, spin for direction)
- \mathbf{P} is generator of translations.
 - Wave momentum shifts coordinate (\mathbf{r}).
 - Intrinsic momentum shifts displacement from equilibrium ($\boldsymbol{\xi}$).

Example: Plane Wave



$$\psi = \sqrt{\frac{\omega Q_0}{2}} \begin{bmatrix} 1 \\ \cos(\omega t - kz) - i \frac{\omega Q_0 k^2}{4\rho c} \sin(\omega t - kz) \\ \cos(\omega t - kz) - i \frac{\omega Q_0 k^2}{4\rho c} \sin(\omega t - kz) \\ 1 \end{bmatrix}$$

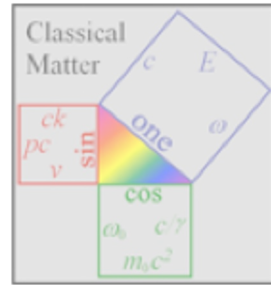
$$s_x = \frac{1}{2} \psi^\dagger \sigma_x \psi = \omega Q_0 \cos(\omega t - kz)$$

$$s_y = \frac{1}{2} \psi^\dagger \sigma_y \psi = 0$$

$$s_z = \frac{1}{2} \psi^\dagger \sigma_z \psi = 0$$

Example: Plane Wave

$$\psi = \sqrt{\frac{\omega Q_0}{2}} \begin{bmatrix} 1 \\ \cos(\omega t - kz) - i \frac{\omega Q_0 k^2}{4\rho c} \sin(\omega t - kz) \\ \cos(\omega t - kz) - i \frac{\omega Q_0 k^2}{4\rho c} \sin(\omega t - kz) \\ 1 \end{bmatrix}$$



$$U_R = \left(-\text{Re}(\psi^\dagger \gamma^5 \boldsymbol{\sigma} \cdot i\nabla \psi) + \frac{1}{4} \check{\omega}_0 \psi^\dagger \gamma^0 \psi \right)$$

$$= \frac{k^2 \omega^2 Q_0^2}{8\rho} \sin^2(\omega t - kz)$$

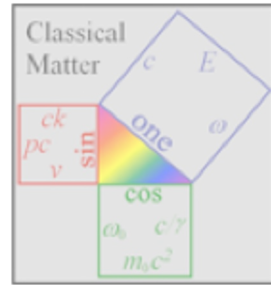
$$K_R = \frac{1}{4} \mathbf{w} \cdot \boldsymbol{\sigma} \psi = \frac{k^2 \omega^2 Q_0^2}{8\rho} \cos^2(\omega t - kz)$$

$$i\partial_t \psi = \left(\gamma^5 \boldsymbol{\sigma} \cdot i\nabla \psi + \frac{1}{2} \check{\omega}_0 \gamma^0 \psi \right) + \frac{1}{2} \mathbf{w} \cdot \boldsymbol{\sigma} \psi$$

$$\mathcal{E} = \mathbf{v} \cdot \mathbf{P} - 2U + 2K_R$$

RHS is NOT Hamiltonian (because nonlinear)²⁰

Conservation Law



Nonlinear Dirac equation:

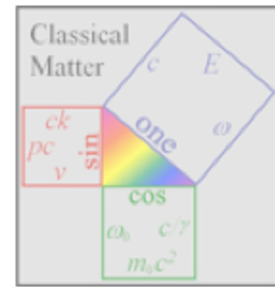
$$\partial_t \psi = -c \gamma^5 \boldsymbol{\sigma} \cdot \nabla \psi - \mathbf{u} \cdot \nabla \psi - \frac{i}{2} (\tilde{w}_0 \gamma^0 \psi + \mathbf{w} \cdot \boldsymbol{\sigma} \psi)$$

Multiply ψ^\dagger and add adjoint:

$$\partial_t (\psi^\dagger \psi) = -\nabla \cdot (\psi^\dagger c \gamma^5 \boldsymbol{\sigma} \psi) - \mathbf{u} \cdot \nabla (\psi^\dagger \psi)$$

Magnitude $\psi^\dagger \psi$ is a conserved quantity.

Exclusion Principle & Potentials



- Wave superposition of “particles” A and B:

$$[\psi_A + \psi_B]^\dagger [\psi_A + \psi_B] = \psi_A^\dagger \psi_A + \psi_B^\dagger \psi_B + \psi_A^\dagger \psi_B + \psi_B^\dagger \psi_A$$

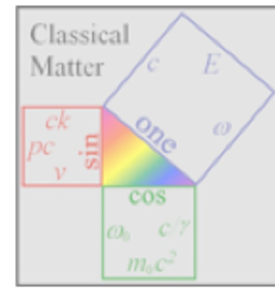
- Interference terms cancel for “independent” particles \Rightarrow exclusion principle:

$$\psi_A^\dagger \psi_B + \psi_B^\dagger \psi_A = 0$$

- Potentials derive from phase shifts introduced to maintain zero interference.

$$\psi_A^\dagger e^{-i\varphi_A} e^{i\varphi_B} \psi_B + c.c. = 0$$

Exclusion Principle & Potentials



Neglect nonlinear terms.

Transformation of momentum operator:

$$\begin{aligned}\psi_A^\dagger (-i\nabla)\psi_A &\rightarrow \psi_A^\dagger e^{-i\varphi_A} (-i\nabla) e^{i\varphi_A} \psi_A \\ &= \psi_A^\dagger (-i\nabla + \nabla\varphi_A)\psi_A\end{aligned}$$

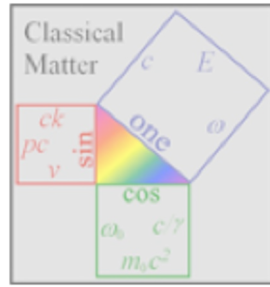
Transformation of Hamiltonian:

$$\begin{aligned}\psi_A^\dagger (\hbar\partial_t + iH)\psi_A &\rightarrow \psi_A^\dagger e^{-i\varphi_A} (\hbar\partial_t + iH) e^{i\varphi_A} \psi_A \\ &= \psi_A^\dagger (\hbar\partial_t + iH + i\hbar\partial_t\varphi_A + i\hbar c\gamma^5 \boldsymbol{\sigma} \cdot \nabla\varphi_A)\psi_A \\ &\equiv \psi_A^\dagger (\hbar\partial_t + iH + ie\Phi - i\gamma^5 \boldsymbol{\sigma} \cdot e\mathbf{A})\psi_A\end{aligned}$$

Electromagnetic potentials:

$$e\mathbf{A} = -\hbar\nabla\varphi_A; \quad e\Phi = \hbar\partial_t\varphi_A$$

Exclusion Principle & Potentials



Change of wave momentum (force): $\frac{d}{dt} P_i =$

$$\psi_A^\dagger \left((\partial_t \partial_i - \partial_i \partial_t) \varphi_A - c \gamma^5 \sigma_j (\partial_i \partial_j - \partial_j \partial_i) \varphi_A \right) \psi_A$$

Multivalued phase \Rightarrow Non-commuting!

Electromagnetic variables:

$$e\mathbf{A} = -\hbar \nabla \varphi_A; \quad e\Phi = \hbar \partial_t \varphi_A$$

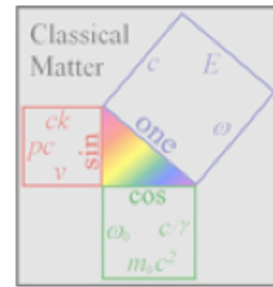
$$\rho_e = e \psi_A^\dagger \psi_A; \quad \mathbf{J} = e \psi_A^\dagger c \gamma^5 \boldsymbol{\sigma} \psi_A$$

$$\mathbf{E} = -\frac{\hbar}{e} \nabla (\partial_t \varphi_A) + \frac{\hbar}{e} \partial_t (\nabla \varphi_A) = -\nabla \Phi - \partial_t \mathbf{A}$$

$$B_k = \epsilon_{kij} \partial_i A_j = -\frac{\hbar}{e} \epsilon_{kij} \partial_i \partial_j \varphi_A$$

$$\frac{d}{dt} \mathbf{P} = \rho_e \mathbf{E} + \mathbf{J} \times \mathbf{B}$$

Exclusion Principle & Potentials



Stokes' Law: $\oint \mathbf{A} \cdot d\boldsymbol{\ell} = \iint (\nabla \times \mathbf{A}) \cdot d\mathbf{S}$

Macroscopic system:

Pure phase shift \Rightarrow magnetic flux quantized:

$$\varphi_A = (m_\phi \phi - \omega t)$$

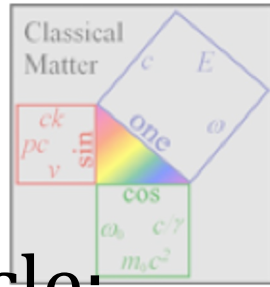
$$\mathbf{A} = -\frac{\hbar}{e} \nabla \varphi_A = -\frac{\hbar m_\phi}{e r \sin \theta}$$

$$\mathbf{B} = \nabla \times \mathbf{A} = \frac{\hbar m_\phi}{e} 2\pi \delta^{(2)}(x, y)$$

$$\oint \mathbf{A} \cdot d\boldsymbol{\ell} = -\frac{\hbar}{e} 2\pi m_\phi \xrightarrow{m_\phi=1/2} \frac{h}{2e}$$

c.f. Hagen Kleinert, *Multivalued Fields in Condensed Matter, Electromagnetism, and Gravitation* (Chapter 4).

Exclusion Principle & Potentials



Radially weighted phase shift from particle:

$$\varphi_A = (m_\phi \phi - \omega t) \left(\frac{e^2}{4\pi\epsilon_0 \hbar c} \right) \frac{c}{\omega r}$$

$$e\mathbf{A} = \left(\frac{e^2}{4\pi\epsilon_0 \hbar c} \right) \frac{c}{\omega r} \left\{ \left(\frac{m_\phi}{r \sin \theta} \right) \hat{\phi} - \frac{m_\phi \phi - \omega t}{r^2} \hat{r} \right\}$$

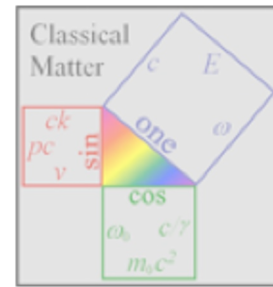
Neglect radial term (no flux). Note $\hbar\omega = m_e c^2$

$$e\mathbf{B} = \left(\frac{e^2}{4\pi\epsilon_0 \hbar c} \right) \frac{c}{\omega} \left(\frac{m_\phi}{r^3 \sin \theta} \right) \hat{\theta}$$

Same flux as e^- dipole moment: $\mu_0 M / 2r$:

$$\oint \mathbf{A} \cdot d\boldsymbol{\ell} = -\frac{m_\phi e}{2\epsilon_0 \omega r} \xrightarrow{m_\phi=1/2} \frac{\mu_0 \hbar e}{4m_e r}$$

Exclusion Principle & Potentials



Radially weighted phase shift:

$$\varphi_A = (m_\phi \phi - \omega t) \left(\frac{e^2}{4\pi\epsilon_0 \hbar c} \right) \frac{c}{\omega r}$$

Phase velocity: $v_\phi = \frac{\omega}{m_\phi} r \sin \theta \rightarrow 2\omega r \sin \theta$

Electric field (Note $\partial_t \varphi_A = -\mathbf{v} \cdot \nabla \varphi_A$):

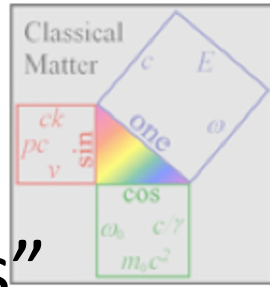
$$eE_i = \hbar(\partial_i v_j) \partial_j \varphi_A + \hbar v_j (\partial_i \partial_j - \partial_j \partial_i) \varphi_A$$

Only first term contributes ($v_\phi = 0$ at $r = 0$)

$$\mathbf{E} = \left(\frac{e}{4\pi\epsilon_0 r^2} \right) \hat{\mathbf{r}}$$

c.f. Herbert Jehle, Phys. Rev. 3(2):306-345, 1971

Matter & Antimatter



Assume vector spherical harmonic “particles”.

Bispinor angular quantum number is half of vector angular quantum number.

Vector: $\ell = 2N + 1 \Rightarrow$ Bispinor: $\ell = \frac{2N+1}{2}$

Odd parity \Rightarrow Distinct mirror image \Rightarrow Fermion

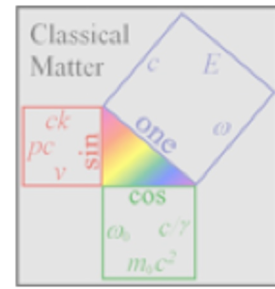
Vector: $\ell = 2N \Rightarrow$ Bispinor: $\ell = N$

Even parity \Rightarrow Same mirror image \Rightarrow Boson

(bosons = antiparticles except $W^{+/-}$)

Assumption valid except for $W^{+/-}$.

Solid Aether Model



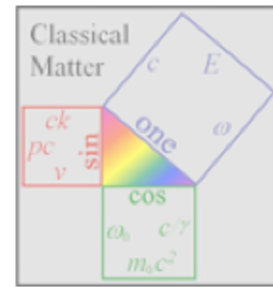
- Lorentz-invariant wave equation (SR).
- “Gravity” is wave refraction (GR).
- Spin is the angular momentum of the medium.
- Mass \Rightarrow spatial localization, velocity rotation.
- Spin & orbital angular momentum.
- Odd/even spherical harmonics: fermions/bosons
- Spatial reflection yields “antiparticles”.
- Exclusion principle and interaction potentials.
- Interpretation of electromagnetism.
- Relation between magnetic flux & electric charge.
- Wave uncertainty principle.

Publications



- Torsion Waves in Three Dimensions: Quantum Mechanics With a Twist," **Found. Phys. Lett.** 15(1):71-83, Feb. 2002.
- "A Classical Dirac Bispinor Equation," in *Ether Space-time & Cosmology, vol. 3*, eds. M. C. Duffy & J. Levy, (Aperion, Montreal, 2009).
- "Exact Description of Rotational Waves in an Elastic Solid," **Adv. App. Clifford Algebras** 21:273-281, 2010.
- *The Wave Basis of Special Relativity* (Verum Versa, 2014)
- "Spin Angular Momentum and the Dirac Equation," **Electr. J. Theor. Phys.** 12(33):43-60, 2015.
- More at: www.ClassicalMatter.org

Related Publications



- Belinfante, F. J. 1939 On the spin angular momentum of mesons. *Physica*, 6:887–898.
- de Broglie LV 1924 *Recherches sur la Theorie des Quanta*, PhD Thesis, (Paris: University of Sorbonne).
- Einstein A 1956 *The Meaning of Relativity* (Princeton: Princeton Univ. Press) Fifth Edition, pp 84-89.
- Evans JC et. al. 2001 Matter waves in a gravitational field: An index of refraction for massive particles in general relativity, *Am. J. Phys.* **69**, 1103--10.
- Gu YQ 1998 Some Properties of the Spinor Soliton, *Advances in Applied Clifford Algebras* **8**(1) 17-29.
- Hestenes D 1990 The Zitterbewegung Interpretation of Quantum Mechanics, *Found. Phys.* **20**(10) 1213-32.
- Jehle H 1971, Relationship of Flux Quantization to Charge Quantization and the Electromagnetic Coupling Constant, *Phys. Rev.* **3**(2):306-345.
- Karlsen BU 1998 Sketch of a Matter Model in an Elastic Universe (<http://home.online.no/~ukarlsen>).
- Kleinert H 1989 *Gauge Fields in Condensed Matter* vol II (Singapore: World Scientific) pp 1259.
- Laughlin R 2005, *A Different Universe* (New York, Basic Books).
- Lee TD and Yang CN 1956 Question of Parity Conservation in Weak Interactions, *Phys. Rev.* **104**, 254.
- Morse PM and Feshbach H 1953a *Methods of Theoretical Physics* vol I (New York: McGraw-Hill Book Co.) pp 304-6.
- Ohanian HC 1986 What is Spin, *Am. J. Phys.* **54**(6):500-5.
- Ranada AF 1983 Classical Nonlinear Dirac Field Models of Extended Particles *Quantum Theory, Groups, Fields, and Particles* ed A O Barut (Amsterdam: Reidel) pp 271-88.
- Rosenfeldr, L. 1940 Sur le tenseur d'impulsion-energie (on the energy-momentum tensor). *Mmoires Acad. Roy. de Belgique*, 18:1–30.
- Rowlands P 1998 The physical consequences of a new version of the Dirac equation *Causality and Locality in Modern Physics and Astronomy: Open Questions and Possible Solutions* (Fundamental Theories of Physics, vol 97) eds G. Hunter, S. Jeffers, and J-P. Vigièr (Dordrecht: Kluwer Academic Publishers) pp 397-402.
- Rowlands P 2005 Removing redundancy in relativistic quantum mechanics *Preprint* arXiv:physics/0507188.
- Schmeltzer I 2012 The standard model fermions as excitations of an ether, in *Horizons in World Physics, vol. 278*, edited by A. Reimer, (Nova Science Publishers).
- Steinberg DJ, Cochran SG, & Guinana MW 1980 A constitutive model for metals applicable at high-strain rate, *J. Appl. Phys.* **51**:1498-1504.
- Takabayashi Y 1957 Relativistic hydrodynamics of the Dirac matter, *Suppl. Prog. Theor. Phys.* **4**(1) 1-80.
- Whittaker E 1951 *A History of the Theories of Aether and Electricity*, (Edinburgh: Thomas Nelson and Sons Ltd.).
- Wilson HA 1921 An Electromagnetic Theory of Gravitation, *Phys. Rev.* **17**: 54-59.
- Yamamoto H 1977 Spinor soliton as an elementary particle, *Prog. Theor. Phys.* **58**(3), 1014--23 .